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The stresses near the free edge of the interface between dissimilar materials under thermal loading can be described by one or two singularity terms and a term which is independent of the distance from the free edge. Some examples are given for the effect of the material properties on the parameters describing the stress distribution.

### INTRODUCTION

If two materials with dissimilar elastic properties and thermal expansion coefficients are bonded together high stresses occur at the free edge of the interface under mechanical loading or after a change in the temperature. In ceramic-metal joints these stresses may cause failure in the ceramic material. The stresses depend on the mentioned physical properties, but also on the size and geometry of the joint and the plastic deformation behaviour of the different components, especially of the solder in a brazed metal-ceramic joint.

In this paper only a joint of two materials is considered and elastic

deformation is assumed.

#### BASIC RELATIONS

In Fig. 1 the general configuration at the free edge is shown, which is characterized by the two angles  $\theta_1$  and  $\theta_2$ . The stresses can be calculated by applying Airy stress function and the boundary and continuity conditions at the free edge and the interface. The stresses near the free edge of the interface can be expressed by the relation

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$$\sigma_{ij}(\mathbf{r},\theta) = \sum_{\mathbf{k}=1}^{N} \frac{K_{\mathbf{k}}}{(\mathbf{r}/\mathbf{L})^{\omega_{\mathbf{k}}}} \cdot f_{ijk}(\theta) + \sigma_{o} f_{ijo}(\theta)$$
(1)

where  $\boldsymbol{L}$  is a characteristic length of the component,  $\omega_k$  are the stress exponents,  $K_k$  stress intensity factors and  $f_{ijk}$  are angular functions. Dependent on the geometry and the elastic constants, one or two  $\omega_{\boldsymbol{k}}$  are positive leading to stress singularities.

For mechanical loading there is  $\sigma_0=0$  (excluding the case of a body with a crack). For thermal loading, however,  $\sigma_0$  is a major contribution to the stress field also near the singularity point. For thermal loading  $\sigma_{\rm 0}$ and  $K_k$  are proportional to  $\Delta\alpha\cdot\Delta T$  and functions of the Dundurs parameters and the angles  $\theta_1$ ,  $\theta_2$ .  $\Delta T$  is the difference between the actual temperature  $T_2$  and the temperature  $T_1$ , where the joint is stress free. Δa is the difference between the effective thermal expansion coeffi-

Signature Streets:
$$\Delta \alpha = \begin{cases} \alpha_1 - \alpha_2 & \text{for plane strees} \\ \alpha_1(1 + v_1) - \alpha_2(1 + v_2) & \text{for plane strain} \end{cases}$$
(2)

The Dundurs parameters are dependent on the elastic constants of the two materials (1).

The  $\omega_{\boldsymbol{k}}$  and the angular functions depend only on the elastic constants and the geometry.  $\sigma_0, \omega_k$  and the  $f_{ijk}$  can be calculated analytically, whereas the Kk have to be determined with numerical methods, for instance the finite element method.

## THE JOINT WITH $\theta_1 = -\theta_2 = 90^{\circ}$

For this joint the  $\sigma_0\text{-term}$  is given in cartesian coordinates by

$$\sigma_{\mathbf{x}} = \tau_{\mathbf{x}\mathbf{y}} = 0$$
  $\sigma_{\mathbf{y}} = \sigma_{\mathbf{0}}$  (3)

with

$$\sigma_{\rm o} = \Delta T \cdot \Delta \alpha \frac{1}{E_1^{*-1} - E_2^{*-1}}$$
 (4)

with 
$$\sigma_o = \Delta T \cdot \Delta \alpha \frac{1}{E_1^{*-1} - E_2^{*-1}}$$
 
$$E_i^* = \frac{E_i}{v_i} \quad \text{for plane stress} \quad \frac{E_i}{v_i \, (1 + v_i)} \quad \text{for plane strain}$$

In Figures 2-4,  $\sigma_0$ ,  $\omega$  and K (determined from FEM) are plotted versus In Figures 2-4,  $\sigma_0$ ,  $\sigma_0$  and  $\sigma_0$  are grained from FEW) are plotted versus,  $(E_1*-E_2*)/(E_1*+E_2*)$ . It can be seen that for  $E_1*\to E_2*$  there is  $\sigma_0\to 0$ ,  $\sigma_0\to \infty$  and  $\sigma_0\to \infty$  and  $\sigma_0\to \infty$  and  $\sigma_0\to \infty$  are found that the region of the section between the ratio  $\sigma_0\to \infty$  $\omega \to 0$ . It was found that there is a unique function between the ratio  $K/\sigma_0$ and the exponent  $\omega$  (Fig. 5) which can be expressed for H/L>1 by (2)

the exponent 
$$\omega$$
 (Fig. 5) which can be exponent  $\omega$  (5)
$$- K/\sigma_0 = 1 - 2.89 \omega + 11.4\omega^2 - 51.9\omega^3 + 135.7\omega^4 - 135.8\omega^5 \qquad (5)$$

Thus it is possible to calculate the stress distribution analytically from eq. (1), without any finite element caculation.

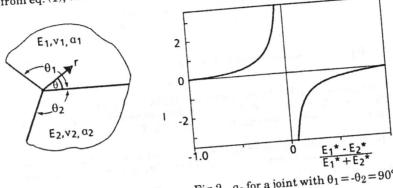


Fig.1 General configuration at the edge of a joint between dissimilar materials

Fig.2  $\sigma_0$  for a joint with  $\theta_1 = -\theta_2 = 90^\circ$ 

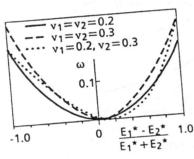


Fig. 3.  $\omega$  versus differences in effective moduli  $(\theta_1 = -\theta_2 = 90^\circ)$ 

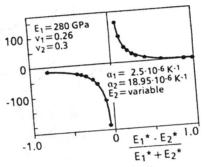


Fig. 4. Effect of difference in effective moduli on  $K_I$  /  $\Delta T$  $(\theta_1 = -\theta_2 = 90^\circ)$ 

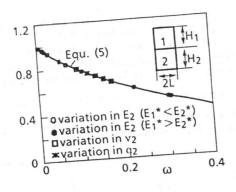


Fig 5.  $K_1/\sigma_0$  versus  $\omega$  for thermal loading

Two conclusions can be obtained from the application of eq. (1).

- 1. There is a large effect of the size of the component on the stresses near the free edge. This can be seen from Fig. 6.
- 2. The stresses are not increasing with increasing  $\boldsymbol{\omega}.$  This can be seen from Fig. 7. Only very close to the free edge the stresses increase with increasing  $\omega$ .

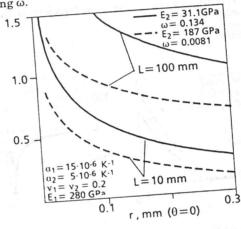


Fig 6. Effect of component size on  $\sigma\theta/\Delta T$  for thermal loading.

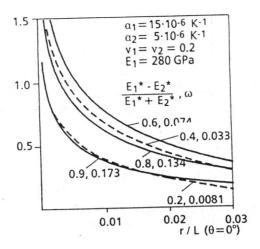


Fig. 7.  $\sigma_{\theta}/\Delta T$  for  $\theta=0^{\circ}$  versus r/L for different  $E_1^*$  -  $E_2^*$ 

## JOINTS WITH ARBITRARY ANGLES $\theta_1$ AND $\theta_2$

As an example a joint with  $\theta_1$  =165° and  $\theta_2$ =-55° is considered. In Fig.8.  $\sigma_0,~K_1,~K_2,~\omega_1,~\omega_2$  are plotted versus  $E_2/E_1$  for  $~E_1=280$  GPa,  $v_1$ =0.26 and  $v_2$ =0.30. It can be seen that

- $\omega_1$  and  $\omega_2$  can be <0 (no singularity) or  $>\!0.$  For  $E_2/E_1>\!30.2$  the  $\omega_k$  are complex and specific considerations are necessary.
- $\sigma_0$  approaches infinity for  $\omega_1 \rightarrow 0$  or  $\omega_2 \rightarrow 0$ .
- $K_1$  approaches infinity for  $\omega_1 \rightarrow 0$  and  $K_2$  approaches infinity for  $\omega_2 \rightarrow 0$ .

In Fig. 9  $K_1$  /  $\sigma_0$  and  $K_2$  /  $\sigma_0$  are plotted versus  $\omega_1$  and  $\omega_2$  respectiveley. It can be seen that  $K_k$  /  $\sigma_0$  = -1 for  $\omega_k$  = 0. Thus the stresses remain finite.

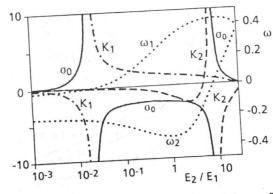


Fig. 8.  $\sigma_0$ ,  $K_1$ ,  $K_2$ ,  $\omega_1$  and  $\omega_2$  versus  $E_2/E_1$  (  $\theta_1=165^\circ$ ,  $\theta_2=-55^\circ$ ,  $E_1=280$  GPa,  $v_1=0,26$ ,  $v_2=0,30$ )

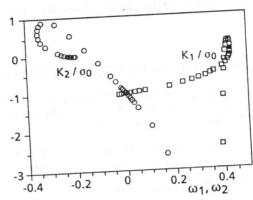


Fig. 9. K/ $\sigma_0$  versus  $\omega$  for combination with  $\theta_1$  = 165°,  $\theta_2$  = -55°

# REFERENCES

- (1) Dundurs, J., J. of Applied Mech., Vol. 36, 1969, p. 650.
- (2) Munz, D. and Yang, Y.Y., J. of Applied Mech., to appear 1992.